Galaxy Zoo: ‘Hanny's Voorwerp’, a quasar light echo?

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ABSTRACT

We report the discovery of an unusual object near the spiral galaxy IC 2497, discovered by visual inspection of the Sloan Digital Sky Survey (SDSS) as part of the Galaxy Zoo project. The object, known as Hanny’s Voorwerp, is bright in the SDSS g band due to unusually strong [OIII]4959, 5007 emission lines. We present the results of the first targeted observations of the object in the optical, ultraviolet and X-ray, which show that the object contains highly ionized gas. Although the line ratios are similar to extended emission-line regions near luminous active galactic nucleus (AGN), the source of this ionization is not apparent. The emission-line properties, and lack of X-ray emission from IC 2497, suggest either a highly obscured AGN with a novel geometry arranged to allow photoionization of the object but not the galaxy’s own circumnuclear gas, or, as we argue, the first detection of a quasar light echo. In this case, either the luminosity of the central source has decreased dramatically or else the obscuration in the system has increased within 10^5 yr. This object may thus represent the first direct probe of quasar history on these time-scales.

Key words: galaxies: active – galaxies: individual: IC 2497 – galaxies: peculiar – quasars: general.

1 INTRODUCTION

The Galaxy Zoo project1 (Lintott et al. 2008) has completed a morphological classification of almost 900 000 objects drawn from the Sloan Digital Sky Survey (SDSS; York 2000; Adelman-McCarthy et al. 2009). By combining classifications made by more than 100 000 participants, it proved possible to compile catalogues of morphology which are of comparable accuracy to those produced

1 www.galaxyzoo.org
by professional astronomers, despite being an order of magnitude larger. The data produced were primarily intended for use in the study of the properties of the population of galaxies (e.g. Bamford et al. 2009), but visual inspection of images from surveys such as the SDSS provides an excellent way of identifying unusual objects within the data set.

In this paper, we discuss an unusual structure, colloquially known as ‘Hanny’s Voorwerp’2 discovered by Hanny van Arkel in the vicinity of the spiral galaxy IC 2497. We report this discovery and present the results of initial follow-up observations in the visible, ultraviolet (UV) and X-ray regions of the spectrum. We consider the emission-line spectrum in detail, and consider possible sources for the observed degree of ionization.

2 PRE-EXISTING OBSERVATIONS OF IC 2497

While there are no pre-existing observations of our target, the neighbouring galaxy IC 2497 is included in several surveys. It has a measured redshift of \( z = 0.050221 \) (Fisher et al. 1995). Assuming, as we will throughout this paper, \( H_0 = 71, \Omega_m = 0.27 \) and \( \Omega_L = 0.73 \) (Dunkley et al. 2009) this redshift corresponds to a luminosity distance of 220.4 Mpc and a scale of 969 pc arcsec\(^{-1}\). With an absolute magnitude of \( M_r = -22.1 \) mag it is a luminous system around 1.7 mag brighter than \( M_r\), (Blanton et al. 2003). The SDSS imaging shows it as a disc galaxy with a large bulge and two fainter spiral arms, as shown in Fig. 1. IC 2497 is also detected at radio wavelengths in the Very Large Array (VLA) Faint Images of the spiral arms, as shown in Fig. 1. IC 2497 is also detected at radio imaging shows it as a disc galaxy with a large bulge and two fainter

Figure 1. The SDSS images of the main galaxy IC 2497 and the Voorwerp. We show each of the five SDSS images (\( u, g, r, i \) and \( z \)) separately as well as a three-colour \( gri \) composite (Lupton et al. 2004). The latter is similar to the Galaxy Zoo image in which the Voorwerp was discovered. IC 2497 is clearly visible in all five bands, while the Voorwerp is strikingly prominent only in the \( g \) band. It is marginally detected in the \( u, r \) and \( i \) bands, and undetected in \( z \). We also indicate the physical scale in \( h_70 \) kpc at the redshift of the system.

1995), with a flux 16.1 ± 0.8 mJy at 1.4 GHz, and hence a radio luminosity of \( L_{\text{1.4 GHz}} = 1.00 \pm 0.05 \times 10^{23} \) WHz\(^{-1}\).

IC 2497 was also detected by the IRAS (Infrared Astronomical Satellite) at 25, 60 and 100 \( \mu \)m, with values from the point-source catalogue giving it an infrared (IR) luminosity of \( L_{\text{IR}} = 3.9 \times 10^{11} L_\odot \) (Sanders & Mirabel 1996), and is thus a luminous infrared galaxy (LIRG). However, inspection of the IRAS data using the Infrared Sky Atlas (IRSA) tool at the Infrared Processing and Analysis Centre (IPAC)\(^3\) web archive shows that the 60 \( \mu \)m measurement (and possibly the others) may be confused with a stronger source about 2 arcmin to its south.

To verify the IRAS fluxes, we used the \textsc{scanpi} web tool from IPAC to retrieve fluxes for each detector crossing of IC 2497, establishing the absence of confusing sources and averaging the scans for measurement. The resulting flux densities were 0.14, 0.22, 2.04 and 3.71 Jy in the 12, 25, 60 and 100 \( \mu \)m bands, respectively, with errors of 0.02, 0.02, 0.02 and 0.06 Jy. Using the far-IR (FIR) parameter from Lonsdale & Helou (1985) the luminosity from 42 to 122 \( \mu \)m is \( 6 \times 10^{44} \text{ erg s}^{-1}\). Despite the high luminosity for such an ordinary-looking galaxy, we note that the FIR energy distribution suggests emission from a source which is colder than most active galactic nucleus (AGN)-dominated sources.

3 IMAGING DATA

3.1 SDSS imaging data and photometric properties

Hanny’s Voorwerp was initially identified in visual inspection of SDSS imaging. In Fig. 1, we present the full SDSS \textit{ugriz} imaging

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2 ‘Voorwerp’ is Dutch for object.
3 NASA/IPAC Extragalactic Database, http://nedwww.ipac.caltech.edu/
4 http://www.ipac.caltech.edu
including a gri colour composite (Lupton et al. 2004) similar to that displayed by the Galaxy Zoo website.

The SDSS photometric data clearly flag the Voorwerp as an unusual object. It is a significant detection only in the g band, where it reaches an apparent magnitude of \( g = 18.84 \). Integrated magnitudes within an aperture of 10 arcsec in the SDSS bands (but in the Pogson logarithmic convention rather than the SDSS sinh style) are \( u = 20.5 \pm 0.15, g = 18.12 \pm 0.08, r = 21.3 \pm 0.1 \) and \( i = 19.8 \pm 0.1 \). The object was not detected in the \( z \) band. Similar distributions between bands are seen at each of the six SDSS photometric objects associated with the Voorwerp, justifying the use of integrated magnitudes. The most unusual is the ‘knot’ to the north-west (NW), which has a different spectral energy distribution from the bulk of the object, being particularly bright in the \( g \) and \( i \) bands. This may suggest contamination by a background source, but a spectrum of the object, being particularly bright in the \( g \)-band image, deeper than that in SDSS, reveals that the Voorwerp is a significantly larger system than was previously apparent in the data shown in Fig. 1. The detected emission extends over \( 18 \times 40 \) arcsec\(^2\) (east–west versus north–south) with additional outlying emission visible to the west.

The SDSS photometric data consist of three 400-s images in each of the \( g, r \) and \( i \) bands (on 2008 January 11) and, on the 2008 January 9, a 600-s image using the wide \( \text{H}\beta \) narrow-band filter [centred on \( \lambda = 4861 \) Å; full width at half-maximum (FWHM) = 170 Å], which at the redshift of IC 2497 traces \( \text{He} \ ii \lambda 4686 \). All four images are shown in Fig. 2. The \( g \)-band image, deeper than that in SDSS, reveals that the Voorwerp is a significantly larger system than was previously apparent in the data shown in Fig. 1. The detected emission extends over \( 18 \times 40 \) arcsec\(^2\) (east–west versus north–south) with additional outlying emission visible to the west.

The morphology of the object is complex, and includes several prominent features. The \( g \)-band images, which are dominated by [O iii] \( \lambda 4959, 5007 \) emission, reveal a lumpy structure, particularly in the part of the object closest to IC 2497. Moving further away, several smaller discrete structures appear that form a nearly round ‘bubble’. This hole is 5.4 arcsec in diameter (corresponding to 4.9 kpc at the distance of IC 2497). The high degree of symmetry seen in this structure poses puzzling questions about its origin.

Remarkably, the Voorwerp is detected in the He\( \alpha \) narrow-band image with the He\( \alpha \) emission in this band coinciding with the brightest features seen in the \( g \) band.

3.2 INT imaging data

We have obtained a series of deeper imaging data from the Wide Field Imager (WFI) at the Isaac Newton Telescope (INT). The data consist of three 400-s images in each of the \( g, r \) and \( i \) bands (on 2008 January 11) and, on the 2008 January 9, a 600-s image using the wide \( \text{H}\beta \) narrow-band filter [centred on \( \lambda = 4861 \) Å; full width at half-maximum (FWHM) = 170 Å], which at the redshift of IC 2497 traces \( \text{He} \ ii \lambda 4686 \). All four images are shown in Fig. 2. The \( g \)-band image, deeper than that in SDSS, reveals that the Voorwerp is a significantly larger system than was previously apparent in the data shown in Fig. 1. The detected emission extends over \( 18 \times 40 \) arcsec\(^2\) (east–west versus north–south) with additional outlying emission visible to the west.

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3.3 He\( \alpha \) imaging data

We also obtained an image of the field through a filter centred on the He\( \alpha \) line, as well as an off-line exposure on the adjacent continuum. These images were taken with the Kitt Peak National Observatory

![Figure 2](https://academic.oup.com/mnras/article-abstract/399/1/129/1086654)
CCD which sampled the image with 0.305 arcsec pixel$^{-1}$ (KPNO 2.1-m telescope on 2008 March 27, using a 2k to the south-west of the IC 2497 nucleus, which is not seen so clearly in any other band. The Voorwerp itself is prominent. We note also the emission source Figure 3.

The characteristic shape of the Voorwerp, including the ‘bub-

ble’ discussed above, is clearly seen. We also note the appearance of a second, resolved source 2.3 arcsec to the west-south-west of the IC 2497 nucleus, suggestive of a double nucleus or a minor merging event. While this is most clearly seen in the Hα image, this source is present in our optical continuum images, showing that it is substantially a continuum object.

### 3.4 Deep continuum imaging in $R$

To provide a better measurement of the red continuum, a total exposure of 110 min in the Bessel $R$ band was obtained on 2008 April 27/28, using the remotely operated 0.9-m telescope of the Southeastern Association for Research in Astronomy (SARA) sited on Kitt Peak. The detector was a 2048 × 2048 pixel E2V chip in an Apogee U42 camera, giving pixel sampling of 0.38 arcsec pixel$^{-1}$. The passband used has Hα and [N ii]λ6583 Å in the red wings of its transmission, so correction for their effects introduces only a small uncertainty. Using the energy zero-points from Fukugita, Shimasaku & Ichikawa (1995) and the same integration region used for total flux from the INT $g$ image, we derive an averaged flux in $R$ across the emitted-wavelength range 5900–6500 Å of $8.8 \pm 1.0 \times 10^{-18}$ erg cm$^{-2}$ s$^{-1}$ Å$^{-1}$.

### 4 SPECTRAL DATA

Spectra covering most of the optical band were obtained with double-spectrograph systems at the 4.2-m William Herschel Telescope (WHT) on La Palma and the 3-m Shane telescope of Lick Observatory. Details of the observations are given in Table 1. The slit width was 2.0 arcsec in both cases, and placement on the sky was nearly identical, passing in both cases through the nucleus of IC 2497, as shown in Fig. 4.

We applied the same reduction procedure to each data set. To eliminate the ripples in sensitivity due to the dichroic beamsplitters in each double spectrograph, which are especially troublesome near [O iii]λ5007 at this redshift, we used the flat-field exposures as obtained, omitting the common step of removing large-scale spectral gradients. After flat-fielding, the spectra thus appeared very blue, but the response curves generated from standard stars were monotonic across almost the entire spectral range and were well fitted

### Table 1. Spectroscopic observations.

<table>
<thead>
<tr>
<th>Telescope</th>
<th>WHT 4.2 m</th>
<th>Lick 3 m</th>
<th>Lick 3 m</th>
</tr>
</thead>
<tbody>
<tr>
<td>Spectrograph</td>
<td>ISIS</td>
<td>Kast</td>
<td>Kast</td>
</tr>
<tr>
<td>Exposure (min)</td>
<td>30</td>
<td>30</td>
<td>2 × 30</td>
</tr>
<tr>
<td>PA (°)</td>
<td>9.5</td>
<td>8.5</td>
<td>8.5</td>
</tr>
<tr>
<td>Slit width (arcsec)</td>
<td>1.97</td>
<td>2.0</td>
<td>2.0</td>
</tr>
<tr>
<td>Airmass</td>
<td>1.49</td>
<td>1.02</td>
<td>1.00</td>
</tr>
<tr>
<td>Dichroic split (Å)</td>
<td>5300</td>
<td>5400</td>
<td>5400</td>
</tr>
<tr>
<td>Blue: wavelength range (Å)</td>
<td>3150–5350</td>
<td>3750–5400</td>
<td>3650–5350</td>
</tr>
<tr>
<td>Dispersion (Å pixel$^{-1}$)</td>
<td>4.88</td>
<td>2.63</td>
<td>2.63</td>
</tr>
<tr>
<td>Spectral FWHM (Å)</td>
<td>12.1</td>
<td>5.4</td>
<td>5.4</td>
</tr>
<tr>
<td>Scale along slit (arcsec pixel$^{-1}$)</td>
<td>0.40</td>
<td>0.72</td>
<td>0.72</td>
</tr>
<tr>
<td>Red: wavelength range (Å)</td>
<td>5160–10 050</td>
<td>5160–7680</td>
<td>5650–7740</td>
</tr>
<tr>
<td>Dispersion (Å pixel$^{-1}$)</td>
<td>5.44</td>
<td>2.33</td>
<td>2.35</td>
</tr>
<tr>
<td>Spectral FWHM (Å)</td>
<td>12.8</td>
<td>7.3</td>
<td>7.3</td>
</tr>
<tr>
<td>Scale along slit (arcsec pixel$^{-1}$)</td>
<td>0.45</td>
<td>0.76</td>
<td>0.76</td>
</tr>
</tbody>
</table>

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by smooth functions. The region containing H\β and [O\III]λλ4959, 5007 falls very close to the rollover wavelength for each dichroic at this redshift, and the derived line ratio is thus very sensitive to how well these transmission ripples can be corrected.

Wavelength calibration was performed using standard lamps at each telescope. For the Lick red spectrum, the Ne lamp lacks lines shortward of 5852 Å, so we supplemented this with λ5577 night-sky emission from object data to constrain the fit further. The blue WHT data have the worst wavelength solution, because the CuAr+CuNe lamp has substantial line blending at low dispersion; the rms scatter of individual line wavelengths about the fit was 0.8 Å, or 0.16 pixels. In the other cases, the line scatter about the adopted fits was 0.11–0.17 Å, or 0.03–0.05 pixels. The line lamps were measured at the beginning or end of the nights, so night-sky lines were used to check for zero-point drifts. In particular, the wavelength scale of the WHT red spectrum requires an offset of about 22 Å. The two-dimensional (2D) spectra (object and standard star) were rebinned to linear wavelength scales, confined to the regions where the wavelength solution was well determined. Nyquist ‘ringing’ occurs at the few per cent level for pixels adjacent to [O\III]λ5007 emission after wavelength rebinning.

Sky subtraction used a third-order Chebyshev function fit to sections of the slit free from significant galaxy light and any obvious emission at the wavelengths corresponding to H\alpha or [O\III]λλ4959, 5007 Å, including a small section between IC 2497 and Hanny’s Voorwerp.

Flux calibration used available standard stars. For the WHT, two standard star observations were used although one was only useful in the red. Three stars were used for the first Lick data set and two on the second Lick night. In this latter case, response curves from the two stars agree well in shape but only at 50 per cent level in intensity. Each of the standard stars has calibrated flux data at 50-Å intervals, except in the deep-red telluric bands, so the sensitivity curves are well constrained; individual flux points scatter about the fit by 0.2 mag. A grey shift was thus introduced to match the mean levels for all observations, reducing this scatter to 0.03 mag.

The merged blue and red WHT spectra are shown in Fig. 6. This represents the flux summed over a region of slit 15–36 arcsec from the nucleus of IC 2497, encompassing the brightest emission from Hanny’s Voorwerp. We use this region in assessing overall spectroscopic properties. Although the Lick spectra are not as sensitive as those obtained with the WHT, they have higher spectral resolution and thus give tighter limits on linewidths. They are crucial in fully resolving the density-sensitive [S\II]λλ6717, 6731 Å doublet.

As a further check, we compare the flux obtained from each of the five spectra where the line could be measured. They give a mean integrated flux of $5.7 \times 10^{-14}$ erg cm$^{-2}$ s$^{-1}$ with rms scatter of 23 per cent. The flux ratio of the λ5007 to λ4959 lines gives an additional check on the errors since the flux ratio should always be 2.93, from statistical weights of the energy levels involved. The measured mean value is 2.92, with rms scatter 10 per cent.

The spectrum is dominated by a series of emission lines (Table 2), with [O\III] at a rest wavelength of 5007 Å by far the most prominent. Using the higher resolution Lick data, comparing with the peak of [O\III]λ5007 emission from IC 2497, and intensity weighting along the slit, we derive a mean intensity-weighted redshift for Hanny’s Voorwerp which is 269 ± 20 km s$^{-1}$ less than that measured for IC 2497. This suggests a genuine physical association between the Voorwerp and IC 2497, rather than a line-of-sight projection effect. The emission spectrum and the accompanying continuum are so dominant that we find only indirect hints of a population of stars within Hanny’s Voorwerp (see Section 6.2).

We can use the SDSS g image to estimate the total [O\III] λ5007 flux from the object for comparison with the small region sampled by the spectrograph slits. We use the energy zero-points for the SDSS system from Fujigita et al. (1995), and incorporate the line wavelengths and equivalent widths from the spectra. The total flux we derive in the λ5007 line is $3.2 \times 10^{-13}$ erg cm$^{-2}$ s$^{-1}$, accounting for a fraction $\sim 0.5$ of the total g intensity. Of this, the deeper INT g image shows that a fraction 0.236 of the intensity of the main body, without outlying patches, falls within the 2-arcsec spectroscopic slit location, so the images give a flux within the spectroscopic slit totalling $4.0 \pm 0.8 \times 10^{-14}$ erg cm$^{-2}$ s$^{-1}$, with the error dominated by the line’s equivalent width against the weak continuum at this redshift, and the derived line ratio is thus very sensitive to how well these transmission ripples can be corrected.

Wavelength calibration was performed using standard lamps at each telescope. For the Lick red spectrum, the Ne lamp lacks lines shortward of 5852 Å, so we supplemented this with λ5577 night-sky emission from object data to constrain the fit further. The blue WHT data have the worst wavelength solution, because the CuAr+CuNe lamp has substantial line blending at low dispersion; the rms scatter of individual line wavelengths about the fit was 0.8 Å, or 0.16 pixels. In the other cases, the line scatter about the adopted fits was 0.11–0.17 Å, or 0.03–0.05 pixels. The line lamps were measured at the beginning or end of the nights, so night-sky lines were used to check for zero-point drifts. In particular, the wavelength scale of the WHT red spectrum requires an offset of about 22 Å. The two-dimensional (2D) spectra (object and standard star) were rebinned to linear wavelength scales, confined to the regions where the wavelength solution was well determined. Nyquist ‘ringing’ occurs at the few per cent level for pixels adjacent to [O\III]λ5007 emission after wavelength rebinning.

Sky subtraction used a third-order Chebyshev function fit to sections of the slit free from significant galaxy light and any obvious emission at the wavelengths corresponding to H\alpha or [O\III]λλ4959, 5007 Å, including a small section between IC 2497 and Hanny’s Voorwerp. The regions labelled 1, 2, 3 and 4 correspond to the ‘zones’ in Table 3.

Figure 4. Slit position for both WHT and Lick data plotted on a Hα image with non-linear scaling in intensity to show detail in both IC 2497 and the Voorwerp. The regions labelled 1, 2, 3 and 4 correspond to the ‘zones’ in Table 3.

### Table 2. Measured emission-line ratios from WHT and Lick spectra. These values are averaged over the ~21-arcsec slice of the Voorwerp summed in Fig. 6, and represent the weighted mean of values from Lick nights 1 and 2 and WHT data (weights 1:2:3). Errors reflect the scatter in the independent measurements when a line was detected in multiple observations, and are otherwise estimated for ([Ne\IV] and [S\II]) from the line intensity and local noise level. The [O\II] λ5007 line has a mean surface brightness of 1.4 $\times 10^{-15}$ erg cm$^{-2}$ s$^{-1}$ within this region.

<table>
<thead>
<tr>
<th>Line</th>
<th>Rest wavelength (Å)</th>
<th>Observed wavelength (Å)</th>
<th>Ratio with H\β</th>
</tr>
</thead>
<tbody>
<tr>
<td>[Ne\IV]</td>
<td>3346</td>
<td>3496</td>
<td>0.2 ± 0.07</td>
</tr>
<tr>
<td>[Ne\II]</td>
<td>3426</td>
<td>3580</td>
<td>0.45 ± 0.07</td>
</tr>
<tr>
<td>[O\II]</td>
<td>3736 + 3729</td>
<td>3897</td>
<td>1.54 ± 0.05</td>
</tr>
<tr>
<td>[Ne\II]</td>
<td>3869</td>
<td>4046</td>
<td>0.83 ± 0.04</td>
</tr>
<tr>
<td>H\γ</td>
<td>3889</td>
<td>4067</td>
<td>0.17 ± 0.05</td>
</tr>
<tr>
<td>[Ne\II]+H\γ</td>
<td>3968 + 3970</td>
<td>4152</td>
<td>0.40 ± 0.03</td>
</tr>
<tr>
<td>H\δ</td>
<td>4101</td>
<td>4294</td>
<td>0.21 ± 0.03</td>
</tr>
<tr>
<td>H\τ</td>
<td>4340</td>
<td>4544</td>
<td>0.48 ± 0.03</td>
</tr>
<tr>
<td>[O\II]</td>
<td>4363</td>
<td>4568</td>
<td>0.12 ± 0.03</td>
</tr>
<tr>
<td>He\i</td>
<td>4686</td>
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<td>0.40 ± 0.02</td>
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<td>4861</td>
<td>5088</td>
<td>1.00</td>
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<tr>
<td>[O\II]</td>
<td>5007</td>
<td>5243</td>
<td>10.5 ± 1</td>
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<tr>
<td>He\i</td>
<td>5876</td>
<td>6154</td>
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<tr>
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<td>0.09 ± 0.02</td>
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<td>6876</td>
<td>3.2 ± 0.3</td>
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<td>[S\II]</td>
<td>9532</td>
<td>9999</td>
<td>2.0 ± 0.3</td>
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</tbody>
</table>
5 SWIFT UV/X-RAY DATA

The Swift satellite was used to obtain UV and X-ray data toward the Voorwerp. Observations with the UV/Optical Telescope (UVOT) and the X-Ray Telescope (XRT) took place for 937 s on the 2008 February 8 and 3816 s on 2008 February 13. The only strong emission line in the Swift UVW2 filter at z = 0.048 might be [C II] at 1909 Å, but emission from this line is strongly weighted to higher density gas such as that found in AGN broad-line emission regions rather than the low-density gas in the Voorwerp (Section 6.2). We also use the XRT on board Swift to search for AGN emission from the larger galaxy or from the Voorwerp itself.

Observations with the UVOT telescope consisted of a total of 4700-s exposure. As shown in Fig. 5 the Voorwerp is a strong UV source, with ~0.3 counts s⁻¹ corresponding, for a flat continuum to a flux of 2.51 × 10⁻¹⁶ erg cm⁻². The UV continuum image shown in Fig. 5 reveals approximately the same structure seen in AGN broad-line emission regions rather than the low-density gas in the Voorwerp (Section 6.2). We also use the XRT on board Swift to search for AGN emission from the larger galaxy or from the Voorwerp itself.

There was no detection of either IC 2497 or the Voorwerp in the X-ray data. Statistics in ‘blank-sky’ regions confirm that fewer than three counts were obtained from either object in 3700 s of integration, giving a count rate of less than 0.001 count s⁻¹. Taking a mean effective area between 2 and 10 keV of 90 cm², the sensitivity of the observations was roughly 7.6 × 10⁻¹⁴ erg cm⁻² s⁻¹. At the distance of IC 2497, this corresponds to a limit of 3.3 × 10³³ erg s⁻¹ between 2 and 10 keV.

6 PHYSICAL CONDITIONS IN THE VOORWERP

6.1 Emission-line ratios and diagnostics

Emission lines provide significant information about the physical conditions in the object and on possible sources of ionization. For the analysis below we concentrate on the spectrum summed across the brightest region (as in Fig. 6).

The density-sensitive [S II] λ 6717/6731 doublet ratio is within the errors, in the low-density limit. Specifically, from the higher dispersion Lick data for which the lines are fully resolved, the ratio is 1.52 ± 0.15; we thus derive an upper limit on the density of $n_e < 50$ cm⁻³.

Detection of the [O III] λ 4363 line provides an estimate of the electron temperature via its ratio with the strong λ 4959, 5007 lines (Peimbert & Costero 1969). The observed ratio corresponds to a range $T_e = 13500 ± 1300$ K.

Evidence for internal reddening from the Balmer decrement is equivocal, with errors in the line ratio which are relatively large for such strong lines because we do not have measurements of Hα and Hβ on the same detector. The ratio Hα/Hβ = 3.2 ± 0.3 corresponds to (foreground screen) reddening $E_{B-V} = 0.12 ± 0.10$ for a Milky Way extinction law, assuming an intrinsic Hα/Hβ ratio of 2.87 (appropriate for a case B recombination and a temperature of 10,000 K; Osterbrock & Ferland 2006). We do not correct our measured value for internal extinction in our discussion; non-zero extinction would increase the luminosity and slightly decrease the ionization parameter derived, and have the net effect of narrowing the bounds we derive on the ionizing luminosity for the central source.

The most unusual feature of the spectrum of the Voorwerp is the presence of strong emission lines associated with high-ionization species such as He II λ 4616 Å and [Ne V] λ 3426 Å. We estimate an ionization parameter $U$ following Penston et al. (1990) and Komossa & Schulz (1997). While the He II/Hβ and [Ne V]/[Ne III] ratios depend on $U$, they also depend strongly on the shape of the ionizing spectrum (Komossa & Schulz 1997). We thus concentrate on the [O III] λ 4363/[O II] λ 3727 ratio, which the models cited find to be more robust. Using an analytical fit to interpolate between models listed by Komossa & Schulz (1997), we find log $U = -2.2$. Together with the electron density, this gives an upper bound on the luminosity of the ionizing source.

Over a wide range of conditions in ionized nebulae, the ratios [N II]/Hα and [S II]/Hα scale broadly with abundances. These are both small in Hanny’s Voorwerp, the [N II] λ 6583 Å line in particular suggesting subsonar abundances (crudely $0.1–0.2 Z_\odot$).

Several diagnostic line ratios show significant changes with position along the slit, in the general sense of ionization increasing southward (away from IC 2497). This is illustrated in Table 3 and Fig. 7. In particular [Ne V]/[Ne III], [O III]/Hβ and He II/Hβ all increase with distance from the nucleus of IC 2497.

6.2 Continuum: recombination, two-photon emission and other sources

Continuum radiation is evident in the spectra, especially in the blue, and the intensity of the Swift UV image suggests that this part of the spectrum is also dominated by the continuum. We consider here its spectral shape and possible constituents. We combine imaging and spectroscopic results, all scaled to encompass the region summed along the slit for the spectrum shown in Fig. 6 (2–arcsec wide,
‘Hanny’s Voorwerp’, a quasar light echo?

Figure 6. Spectrum of Hanny’s Voorwerp obtained with the WHT, summed over the slit section 15–36 arcsec south of the nucleus of IC 2497. The prominent [O III]4959, 5007 lines dominate the detected emission, while the presence of [Ne v] and [He II] lines indicates that the gas is more highly ionized than can be accounted for by starlight. In order to display the fainter lines, the brightest [OIII] line is truncated. Blue and red sections of the spectrum have been merged by resampling to a common wavelength scale, and blended with smoothly varying weights across the range of overlap.

Table 3. Emission-line ratios for four averaged positions across Hanny’s Voorwerp and for the nucleus of IC 2497. The regions used are indicated in Fig. 4. Except for [S II] where we give both lines, where appropriate we refer to the stronger line of a pair so that [O III] represents 5007 Å, [N II] 6583 Å and [O I] 6300 Å. [Ne III] and [Ne V] are the single lines at 3969 and 3426 Å, respectively. The error bars given in parentheses for the nucleus indicate the difference expected from subtracting a plausible range of stellar populations, which is significant for H β because of the relatively strong and uncertain correction for underlying absorption.

<table>
<thead>
<tr>
<th>Zone</th>
<th>Distance (kpc)</th>
<th>[N II]/Hα</th>
<th>[S II] 6717 Å/Hα</th>
<th>[S II] 6731 Å/Hα</th>
<th>[S III]/Hα</th>
<th>[O III]/Hβ</th>
<th>He II/Hβ</th>
<th>[Ne V]/[Ne III]</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>13–16</td>
<td>0.31</td>
<td>0.15</td>
<td>0.08</td>
<td>0.16</td>
<td>9.7</td>
<td>0.34</td>
<td>–</td>
</tr>
<tr>
<td>2</td>
<td>16–19</td>
<td>0.25</td>
<td>0.12</td>
<td>0.08</td>
<td>0.39</td>
<td>10.0</td>
<td>0.34</td>
<td>0.79</td>
</tr>
<tr>
<td>3</td>
<td>19–22</td>
<td>0.15</td>
<td>0.07</td>
<td>0.06</td>
<td>0.58</td>
<td>9.7</td>
<td>0.42</td>
<td>0.65</td>
</tr>
<tr>
<td>4</td>
<td>22–31</td>
<td>0.09</td>
<td>0.07</td>
<td>0.04</td>
<td>0.64</td>
<td>10.7</td>
<td>0.46</td>
<td>0.39</td>
</tr>
<tr>
<td>Nucleus</td>
<td>0.8</td>
<td>1.15</td>
<td>0.27</td>
<td>0.27</td>
<td>–</td>
<td>3.6(1.0)</td>
<td>–</td>
<td>–</td>
</tr>
</tbody>
</table>

15–36 arcsec south of the nucleus of IC 2497 along position angle (PA) = 8°. For the spectroscopic points for both WHT and Lick data means in windows free of strong emission lines were used with errors obtained by combining the internal error of the mean with an external 10 per cent flux-scale error. From images, we use the continuum λ6573 image from the KPNO 2.1-m telescope and the longer exposure in Bessell R (which excludes Hα at this redshift) from the SARA 0.9-m. We also include the Swift UVOT measurement. The UVOT passband includes [C III]λ1909, so we assign error bars reflecting the range of [C III]:[O III] λ5007 ratios seen in ionization cones from Seyfert galaxies with similar ionization levels (Evans et al. 1999).

The equivalent width of Hβ is 360 ± 20 Å in the emitted frame. This means that the continuum contributions from recombination and two-photon decay from the metastable 2S1/2 state of H I are not negligible. The equivalent width of Hβ against the recombination (free–free plus bound–free) continuum is 1350 Å at the derived electron temperature (De Robertis & Osterbrock 1986), so that slightly more than a quarter of the observed continuum near Hβ comes from the plasma. We evaluate these contributions using the analytical expressions from Ferland (1980), Nussbaumer & Schmutz (1984) and Osterbrock & Ferland (2006). We assume a helium abundance of 0.08 by number, and neglect He++. The two-photon continuum is scaled to conform to the low-density limit with no collisional de-excitation. The sharp Balmer jump is smoothed in practice by the pseudo-continuum produced by the confluence of high-order Balmer emission lines, which blends smoothly into the Balmer continuum. We have approximated this effect in Fig. 8 based on spectrophotometry of the planetary nebula Jonckheere 900 and the Seyfert galaxy NGC 4151, obtained using the 2.1-m telescope on Kitt Peak (Keel 1987). These objects bracket the linewidths seen in Hanny’s Voorwerp; we have logarithmically interpolated the pseudo-continuum in linewidth, obtaining a shape which is roughly linear in flux between 3646 and 3927 Å.

As shown in Fig. 8, the nebular continuum is a significant fraction of the total in this object and most of the emission just shortward
of the Balmer jump can be attributed to it. However, the optical emission, roughly flat in $F_\lambda$ longward of 5500 Å, must come from processes other than those that produce the continuum. Since the normalization of the free–free continuum is set directly from the H$\beta$ equivalent width with error $\pm$6 per cent, this excess continuum in the optical data is detected at high significance. Most of the flux measured near 2000 Å is also greatly in excess of the two-photon continuum and comes from other sources. The longer wavelength continuum could plausibly represent direct starlight. However, the slope of this residual continuum from the optical to mid-UV is very steep (roughly $\lambda^{-3}$), which may suggest a scattered-light component. Such a contribution should have a strong polarization signature. Given the errors, it is not clear whether we detect any excess

6.3 Velocity structure

Significant velocity structure appears in several emission lines, and in both sets of spectral data. This is best shown in [OIII] $\lambda$5007 in the Lick data, which has higher spectral resolution than the WHT spectrum. Fig. 9 shows the velocity offset from [OIII] $\lambda$5007 in the nucleus of IC2497, compared to that in the Voorwerp. Errors were estimated following Keel (1996), with a floor of 10 km s$^{-1}$ (corresponding to 0.07 pixel) from pixel centroiding. The peak-to-peak amplitude of this velocity slice is about 90 km s$^{-1}$. Our slit location samples the edge of the ‘hole’ which is a prominent feature in our images; the intensity peaks at 23 and 27 arcsec seen in Fig. 9 lie along its rim. It may be significant that this region has the most negative radial velocities we observe.

7 PHYSICAL CONDITIONS IN THE NUCLEUS OF IC 2497

Spectral lines were also detected toward the nucleus of IC 2497. Of particular interest are H$\alpha$, [O I]$\lambda\lambda$6300 Å, [OIII]$\lambda\lambda$3736, 3729 Å and [OIII]$\lambda\lambda$4959, 5007 Å, from which detections we are able to confirm that the galaxy is in fact a Low-Ionization Nuclear Emission-Line Region (LINER) AGN with, taking into account the underlying stellar absorption, [N II]/H$\alpha = 0.8 \pm 0.15$ (Heckman 1980). The [OIII]$\lambda$3727/[OIII]$\lambda$5007 ratio is $1.33 \pm 0.16$. Other important ratios are given in Table 3.

Detection of [S II]$\lambda\lambda$6717, 6731 allows us to calculate the ionizing flux. An ionization parameter for the circumnuclear gas of $10^{-1.2}$ is found. From the measured [S II]$\lambda$6717/[S II]$\lambda$6731 ratio of
1.02 ± 0.05 we derive an electron density of ~560 ± 150 cm⁻³ for the centre of IC 2497, similar to values found for the narrow-line region in Seyfert galaxies (Peterson 2003; Bennert et al. 2006; Osterbrock & Ferland 2006). As often happens in analysis of LINERs, correction for the underlying starlight is a particular issue for Hβ. Most of the stellar features are well fitted by a template based on elliptical galaxies (Kennicutt 1998), but an absorption blend around 3850 Å is deeper in the template than in IC 2497. This may hint at a population of younger stars. The effect of such a population would be to simultaneously decrease the [O III]/Hβ ratio, decrease the reddening needed to account for the He II/Hβ ratio and increase the implied ionization level if the [O III]/[O II] ratio were to be corrected by an appropriate amount to match the He II/Hβ ratio. In view of these uncertainties, we assign a relatively large error to the [O III]/Hβ ratio and do not attempt to correct for reddening. Such a correction would, in any case, reduce the derived ionization parameter, so this is a conservative approach in evaluating the level of nuclear activity ionizing the gas.

The line ratios in both IC 2497 and Hanny’s Voorwerp are illustrated in the ‘BPT diagram’ (Baldwin, Phillips & Terlevich 1981) shown in Fig. 10. The trend to increasing ionization with greater distance from IC 2497 is clearly seen in the data for the Voorwerp, although all the points fall in the part of the BPT diagram defined by the Seyfert regime, whereas IC 2497 lies near the boundary between the parts of this diagnostic diagram associated with LINERs and Seyfert nuclei.

8 DISCUSSION

8.1 Ionizing the Voorwerp: photoionization versus shocks

Our observations suggest that Hanny’s Voorwerp is a low-density gas-rich object, illuminated by a hard ionizing radiation field impinging on the gas. The source of the gas may be IC 2497 itself, or the Voorwerp may be an independent dwarf galaxy. This latter case is suggested by the low derived metallicity, similar to those found for dwarf galaxies by Tremonti et al. (2004).

Gas can be highly ionized either through photoionization by a continuum extending to high energies (soft X-rays in this case, since Ne⁴⁺ has an ionization threshold near 100 eV) or fast shocks. The shock interpretation is difficult to sustain in this instance, for several reasons. Shock velocities of 400 km s⁻¹ are needed to produce strong He II and [Ne V] emission (Dopita & Sutherland 1996), and such velocities are far beyond the radial velocity range of 90 km s⁻¹ observed here. The lack of a systematic correlation between either extreme velocities or velocity gradients and [O III] 5007 Å surface brightness (Fig. 9) argues against large-scale shocks as the means of energy input. Finally, shock models give relations among electron temperature, as measured via the [O III] λ 5007/λ 4363 ratio, and ionization indicators such as He II/Hβ, which require much higher electron temperatures than we see in this case (Evans et al. 1999), typically $T_e \approx 2 \times 10^4$ K.

Although imaging of the Voorwerp at a wide range of wavelengths (including UV imaging and g, r and i bands) reveals the presence of a bubble-like structure which is ~5 kpc across, and might represent a kind of expanding Strömgren sphere, powered by a heavily obscured central source, nothing in the available data suggests such a source. Instead, we must look for a source of ionization external to the Voorwerp itself. It has a similar redshift to IC 2497, suggesting a genuine physical association. Moreover, the increase in ionization level observed across the Voorwerp, decreasing with distance from IC 2497 supports the hypothesis that the neighbouring galaxy is the direct or indirect source of the ionization.

One possible counterpart to the Voorwerp which is the result of the action of a jet is Minkowski’s object (MO), a blue object near NGC 541 within galaxy cluster Abell 194 (Minkowski 1958; van Breugel et al. 1985; Croft et al. 2006). There is strong evidence that star formation observed in MO was triggered by a radio jet from NGC 541; and we can thus compare this exotic object with Hanny’s Voorwerp to look for evidence of a similar origin. Without a detailed search for such a jet in the IC 2497 system it is difficult to say for certain, but there are important observed differences between MO and the Voorwerp. In particular, optical emission from MO is dominated by [O III] and Hα, whereas in the Voorwerp both of these lines are much weaker than the main [O III] 4959, 5007 line. MO also exhibits bright continuum emission, whereas the emission lines are clearly dominant in the Voorwerp spectrum. These results suggest that the source of the Voorwerp’s ionization is different from that in MO; not hot stars, but something else.

It is also unlikely that the energy input results from direct interaction with outflows from IC 2497, such as radio jets. Jozsa (2009) report the detection of such a jet associated with the galaxy, but as noted above, shocks from such an interaction would also have to be much faster than the observed velocity range of the gas to account for the high level of ionization.
8.2 An AGN in IC 2497?

Having ruled out shocks and interaction with radio jets as the cause of the ionization of the Voorwerp, we next consider a possible AGN in IC 2497. This hypothesis is supported by the observed strength of high-ionization species such as He II and [Ne v], which distinguish this object from typical star-forming regions. The best match to these emission-line ratios (as seen in Fig. 10) occurs for gas under conditions similar to those seen in the narrow-line regions of AGN (Leipski et al. 2007; McCarthy 1993), particularly the distant gas forming the ‘extended emission-line regions’ tens of kiloparsecs in size seen around some quasi-stellar objects (QSOs) and radio galaxies (see summaries by Fu & Stockton 2009; Stockton, Fu & Canalizo 2008), with typical [O iii] λ5007 A luminosity exceeding $10^{42}$ erg s$^{-1}$. They are most prevalent accompanying radio-loud quasars but are not structurally related to either the radio sources or host galaxies.

We now constrain the strength of any AGN in several ways: obtaining an upper limit from the lack of an X-ray detection, both upper and lower bounds from the observed emission-line spectrum, and the level of possibly absorbed AGN radiation from the IRAS observations discussed in Section 2.

A lower limit to the required energy input to the gas comes from straightforward energy balance – the number of ionizations and recombinations must match, and the rate of emission of ionizing photons must be at least sufficient to power the observed recombination lines. The integrated Hβ luminosity of the Voorwerp is $1.4 \times 10^{41}$ erg s$^{-1}$. For typical nebular conditions and a flat ionizing continuum, one in 12.2 recombinations cascades through the Hβ transition and one in 9.1 for Hα (table 4.4 in Osterbrock & Ferland 2006). The fraction of the ionizing luminosity (between H and He ionization edges) reprocessed into line emission depends on both the optical depth (making the derived luminosity a lower limit) and covering fraction. In our deepest g image, the emission subtends approximately 38\° about the nucleus of IC 2497, which would correspond to a covering fraction of ~0.03 if it is comparably deep along the line of sight. For a flat ionizing continuum ($F_\nu \propto \nu^{-1}$), this gives a required ionizing luminosity $>1.0 \times 10^{45}$ erg s$^{-1}$; the X-ray luminosity is comparable for this continuum slope.

Since we have an upper limit to the electron density, we can use the ionization parameter in the gas to provide an upper limit to the incident continuum flux and hence luminosity. For ionization parameter $U = 0.006$ and $n_e < 50$ cm$^{-3}$, the local density of ionizing photons will be $<0.32$ cm$^{-3}$. Using the mean projected separation of the Voorwerp from the core of IC 2497, 20 kpc, as the distance the ionizing source must have an isotropic output of $Q_{\text{ion}} < 9.5 \times 10^{36}$ s$^{-1}$. For a flat continuum shape, this corresponds to $L_{\text{ion}} < 3.2 \times 10^{45}$ erg s$^{-1}$, and again a comparable X-ray output would be expected. These two emission-line arguments thus bound the required ionizing luminosity in the range 1–3 $\times 10^{45}$ erg s$^{-1}$.

We can place a limit on any nuclear activity in IC 2497 with the Swift X-ray data, described in Section 5. Assuming an unabsorbed, AGN-like ($\nu^{-1}$) spectrum between 2 and 10 keV, these data rule out relevant AGN luminosities; for a flat, unabsorbed continuum the limit derived from our Swift observations is more than 3 dex below the required luminosities. This conclusion holds unless, along our line of sight, most of the flux up to 5 keV is absorbed (more specifically, using the xspec web tool at the High Energy Astrophysics Science Archive Research Center (HEASARC) site, this means equivalent $N_H < 10^{24}$ cm$^{-2}$).

The key issue is therefore whether IC 2497 could host a powerful AGN which remains active and luminous, but is so deeply obscured as to elude our observations so far.

The FIR data from the IRAS survey (see Section 2) suggest a FIR luminosity of $1.5 \times 10^{44}$ erg s$^{-1}$, an order of magnitude less than the energy requirements we find for ionizing luminosity. This observed value, which applies to the integrated flux of the galaxy, includes any contribution from the disc of what is a very luminous spiral on top of the AGN component. In comparing the total FIR output to the isotropic ionizing flux, we implicitly assume a geometry in which most of the radiation is intercepted by some thick, roughly toroidal structure, suggested by known AGN in which strong obscuration shapes the spatial extent of the escaping radiation (Tadhunter et al. 1999; Mason et al. 2006). Other structures are possible; if the obscuring material is a thinner annulus seen edge-on, or patchy and of small covering fraction but blocking our line of sight, the amount of energy absorbed could be proportionally smaller. Such material would have to satisfy the column density constraints from X-ray observations.

However, this simple picture is unlikely to apply here as the nuclear gas in IC 2497 sees very little ionizing radiation. The emission spectrum from the nucleus of IC 2497 is quite representative of the LINERs often found in early-type spirals, and well explained by a power-law continuum of low-ionization parameter $\log U \approx -3.5$. Obscuration strong enough to block our line of sight and thus hide the core seems unlikely without raising the ionization level of circumnuclear material beyond the observed value. In addition, the ionization cones seen in galaxies such as Cygnus A and NGC 1068 are smaller scale features which surround the nucleus rather than illuminating distant patches of ionization such as we see here.

The region of gas which is seen in emission in a LINER is typically within 0.5 kpc from the nucleus (Preito, Maciejewski & Reunanen 2005), consistent with the region included in the seeing disc of the nucleus for our spectra. If this distance is an appropriate estimate for conditions in IC 2497, then gas which is ~40 times closer to the central ionizing source than the Voorwerp must be seeing an ionizing flux less than half of that seen by the more distant gas. In order to reconcile these observations, we would be forced to postulate some geometry which allows ionizing radiation to escape from the galaxy only through a channel some 20\° in half-angle, without encountering a significant density of the galactic interstellar medium.

Furthermore, the Swift/XRT observation also rules out a luminous, Compton-thick AGN currently residing in IC 2497. Consider a heavily obscured AGN with an intrinsic column density of $1 \times 10^{24}$ cm$^{-2}$, i.e. Compton thick, and an unabsorbed luminosity of $10^{45}$ erg s$^{-1}$ which matches the ionizing luminosity required to explain the ionization in Hanny’s Voorwerp. While most of intrinsic flux is absorbed, Levenson et al. (2006) find that approximately 1 per cent of the continuum’s intrinsic flux is detected in reflection for seven Compton-thick AGN studied with the Chandra X-Ray Observatory. If IC 2497 hosts a $10 \times 10^{45}$ erg s$^{-1}$ Compton-thick AGN, the expected observed 0.2–10keV flux is $1.72 \times 10^{-12}$ erg s$^{-1}$ cm$^{-2}$ and Hα would predict 130 Swift/XRT photons in 3700 s. Therefore, either the reflection fraction in IC 2497 is much smaller than that indicated by Levenson et al. (2006) or IC 2497 does not currently host a sufficiently powerful heavily obscured AGN. What these data alone do not rule out is either a moderately obscured AGN in IC 2497 or even something similar to NGC 1068 which is often regarded as the prototypical example of a Seyfert 2 nucleus. Harder X-ray observations will be necessary to rule out this latter case.

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Thus, from our knowledge of the properties of IC 2497, using observations from the IR through to the X-ray, it is difficult to identify a present-day AGN as the source for the high levels of ionization seen in the Voorwerp, and this leads us to consider an alternative hypothesis.

8.3 A quasar light echo?

In the absence of an ionizing source, we conclude that the Voorwerp was ionized by a source which is no longer active. We hypothesise that IC 2497 underwent an outburst, reaching quasar luminosities, and that we see material which lies close to the light-echo (or constant time-delay) ellipsoid (Couderc 1939) and is illuminated and ionized by this prior outburst.

The first astronomical detection of a light echo, around Nova Persei 1901, was described by Kapteyn (1902). This discovery has been followed by the discovery of simple scattering echoes from – most famously – SN 1987A, the eruptive variable V838 Monocerotis (Bond et al. 2003), and from more distant extragalactic supernovae (e.g. Rest et al. 2008a). Light echoes have recently been exploited to measure the spectra of historical supernovae, and deduce their spectroscopic classifications (Rest et al. 2008b). If our hypothesis is correct, the Voorwerp represents the first detection of the phenomenon with a source that lies on galactic rather than stellar scales.

The separation of the Voorwerp from IC 2497 is between 45 000 and 70 000 light yr, depending on the angle of projection. For a true light echo, as grains are forward scattering, the most favourable scattering geometries for UV dust reflection will place the Voorwerp in front of IC 2497. This suggests that an outburst, or perhaps the end of a longer luminous phase, must have taken place ~10^8 yr ago (referred to the epoch at which we observe IC 2497). The use of ‘light echo’ would be fully consistent with previous usage only for the dust-scattered component which we infer for the UV continuum. The recombination time-scale at the low densities we measure is >8000 yr, but not trivial compared to the light-travel times involved, so the observed emission-line response (‘photoionization echo’) would be more spread in depth than would be the case for pure reflection.

It has long been clear that the AGN population evolves over time (see e.g. Boyle et al. 2000; Wolf et al. 2003; Richards et al. 2006), but it is harder to constrain the time-scales on which individual objects undergo change. The connection between AGN and mergers suggests that the subsequently triggered AGN episodes last typically for 10^9 yr (Stockton 1982; Bahcall et al. 1997) and may last up to 10^9 yr (Bennert et al. 2008). The presence of young stellar populations in many quasar host galaxies suggests that their activity is connected to starbursts with a similar time-scale of ~10^6 yr (Canalizo & Stockton 2001; Miller & Sheinis 2003). At the other end of the scale, there have been numerous detections of AGN which flare on time-scales of years (Storchi-Bergmann, Baldwin & Wilson 1993; Cappellari et al 1999). The time-scale we infer for the shutdown of activity in IC 2497, of ~10^6 yr, is intermediate between these extremes. Short time-scales ~10^6 yr have been suggested for episodes of luminous AGN activity both from the distribution of derived Eddington ratios (Hopkins & Hernquist 2009) and statistics of QSO absorption systems at high redshift (Kirkman & Tytler 2008).

The lowest redshift quasar in the SDSS Data Release 5 (DR5) catalogue (Schneider et al. 2007) lies at z = 0.08, but this sample systematically excludes systems at lower redshift. Our best comparison is with Barger et al. (2005). Taking the 2–8 keV luminosity of 10^{44} (a conservative estimate for the flux required to produce the ionization fraction we observe) the local space density of such luminous AGN is no greater than 3 x 10^{-7} Mpc^3. This suggests that there should be one such system at a redshift of z < 0.04, so the presence of such activity in IC 2497, while unusual, is not entirely unexpected.

If the obscuration along the line of sight to the Voorwerp has remained constant, then the AGN in IC 2497 must have undergone either a very bright flare or else reached the end of an extended period of high luminosity. In either case, detailed observation of the Voorwerp would enable us to reconstruct the history of the source, probing AGN variability on time-scales of 10^6 yr for the first time. This hypothesis suggests further observations which could test it, and, if it is correct, uncover the details of the object’s history. We would expect the scattered continuum to be polarized and show broad QSO emission lines in reflection; this spectral signature would be brightest in the UV, possibly within the range of Galaxy Evolution Explorer (GALEX) for such a large and diffuse target. The variation in ionization parameter might trace changes in the ionizing luminosity; measurements of the density across the object could separate density and time effects. The origin of the gas (and scattering dust) in the Voorwerp may have been a dwarf galaxy, probably close enough to IC 2497 to have been tidally disrupted. Near-IR imagery at high resolution may be the best way to search for star clusters from the pre-existing stellar population with minimal interference from the very blue scattered light and the nebular continuum emission.

9 CONCLUSION

We have presented observations of Hanny’s Voorwerp, an object first identified through visual inspection of the SDSS as part of the Galaxy Zoo project. The object, near to and at the same redshift as IC 2497, a spiral galaxy, is highly ionized and has a spectrum dominated by emission lines, particularly [O II]λλ4959, 5007, with no sign of any contribution from a stellar component to the Voorwerp itself. Both the Voorwerp and its neighbouring galaxy are strong UV sources, but neither was detected in X-ray observations carried out with the Swift satellite. This lack of X-ray detections, and the limits derived from IRAS observations of IC 2497, provides a strong constraint on the luminosity of any ionizing source. We are left with two possible conclusions. Either an AGN in IC 2497 is heavily obscured but still able to ionize the Voorwerp, which extends over almost 20°, or else the ionizing source is no longer present. In the latter case, the Voorwerp represents the first instance of a light echo being seen from a quasar-luminous AGN. In either case IC 2497 furnishes a nearby example of a galaxy which either is, or was shortly before the epoch at which we observe it, a quasar host galaxy.

Detailed further observations, particularly observations in the radio and deep optical imaging, will be required to confirm our hypothesis. However, it is clear that such a light echo would provide an unusual – possibly unique – opportunity to probe the variation of an AGN on time-scales of ~10^8 yr, reconstructing its history by observing echoes from different parts of the Voorwerp.

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