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# Deep spectroscopy of the $M_V \sim -14.8$ host galaxy of a tidal disruption flare in A1795\*

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## ABSTRACT

A likely tidal disruption of a star by the intermediate-mass black hole (IMBH) of a dwarf galaxy was recently identified in association with Abell 1795. Without deep spectroscopy for this very faint object, however, the possibility of a more massive background galaxy or even a disc-instability flare from a weak active galactic nucleus (AGN) could not be dismissed. We have now obtained 8 h of Gemini spectroscopy which unambiguously demonstrates that the host galaxy is indeed an extremely low mass ( $M_* \sim 3 \times 10^8 M_\odot$ ) galaxy in Abell 1795, comparable to the least massive galaxies determined to host IMBHs via other studies. We find that the spectrum is consistent with the X-ray flare being due to a tidal disruption event rather than an AGN flare. We also set improved limits on the black hole mass ( $\log[M_\bullet/M_\odot] \sim 5.3\text{--}5.7$ ) and infer a 15 yr X-ray variability of a factor of  $\gtrsim 10^4$ . The confirmation of this galaxy–black hole system provides a glimpse into a population of galaxies that is otherwise difficult to study, due to the galaxies' low masses and intrinsic faintness, but which may be important contributors to the tidal disruption rate.

**Key words:** galaxies: clusters: individual: Abell 1795 – galaxies: distances and redshifts – galaxies: dwarf – galaxies: kinematics and dynamics – galaxies: nuclei – X-rays: bursts.

## 1 INTRODUCTION

Massive black holes (MBHs; mass  $M_\bullet \gtrsim 10^5 M_\odot$ ) are thought to occur in most if not all galaxies and are thought to be an integral component of galactic formation and evolution. MBHs become progressively more difficult to observe at lower masses, however, due

to their diminishing influence on the surrounding stellar population and weaker accretion capability. In addition, the broad emission-line region of active galactic nuclei (AGNs) may be strongly regulated for  $M_\bullet \lesssim 10^6 M_\odot$  (Chakravorty, Elvis & Ferland 2014). Such intermediate-mass black holes (IMBHs;  $10^2\text{--}10^6 M_\odot$ ) are of great interest for the purpose of studying the physical nature of these relationships, and for devising models which can successfully predict the low-mass and high-mass tails of the black hole distribution (e.g. Volonteri & Natarajan 2009). In particular, the presence of IMBHs in the dwarf galaxies of galaxy clusters is an issue of interest to models of galaxy formation and evolution, as these galaxies numerically dominate the overall galaxy population.

The disruption of a star by such an MBH may temporarily produce a long-lived ( $\sim$ months–years) X-ray flare (Rees 1988; Ulmer 1999) sufficiently luminous ( $L_X \gtrsim 10^{42}$  erg s<sup>−1</sup>) to distinguish the

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emission from that of accreting stellar-mass black holes. Such tidal disruption events (TDEs) in dwarf galaxies therefore provide an important complement to studies of low-mass AGNs (Reines, Greene & Geha 2013) and weakly accreting IMBHs in low-mass galaxies (Miller et al. 2012), and hold implications for dynamical and evolutionary studies of galactic nuclei (e.g. Wang & Merritt 2004), including the rate of gravitational wave events from the disruption of white dwarfs by IMBHs (Sesana et al. 2008).

Maksym et al. (2013) recently identified a likely tidal disruption flare in a faint ( $V = 22.46$ ) galaxy associated with Abell 1795, WINGS J134849.88 + 263557.5 (hereafter WINGS J1348) as part of a survey to find TDEs in archival X-ray observations of galaxy clusters (Maksym, Ulmer & Eracleous 2010). Donato et al. (2014) also later independently identified this flare and came to similar conclusions as to its origins. These studies argued that WINGS J1348 was likely a member of Abell 1795 with a projected distance of only  $\sim 50$  kpc from the cluster core, and that an earlier *Extreme Ultraviolet Explorer* (*EUVE*; Bowyer & Malina 1991) flare of indeterminate nature (identified by Bowyer, Berghöfer & Korpela 1999) had the same origins. Such conclusions implied that the host galaxy was probably a dwarf galaxy ( $M_V \sim -14.8$ ,  $r \sim 300$  pc) hosting a possible IMBH ( $M_\bullet \sim 3 \times 10^5 M_\odot$ , from assumed Eddington accretion and from the  $M_\bullet - L_{\text{bulge}}$  relationship). If so, it would be one of the least massive galaxies confirmed to host an IMBH, comparable to the most extreme examples from Reines et al. (2013) and Miller et al. (2012).

The conclusions of Maksym et al. (2013) and Donato et al. (2014) rely directly on redshifts derived from template fitting of galaxy photometry and the likelihood of an associated galaxy being a true cluster member. Both studies analysed optical spectra which placed limits on any putative prior AGN which could produce an extreme X-ray flare. But in the absence of such emission lines, neither spectrum was sufficiently deep to obtain a spectral redshift via absorption lines, given the faintness of the galaxy. Seyfert 2 galaxies have been observed to commonly produce extreme long-term variability (Saxton et al. 2011), so WINGS J1348 could in principle be a flaring Seyfert 2 at  $z \gtrsim 0.4$  with weak emission lines (Maksym et al. 2013). Alternately, even if WINGS J1348 has been quiescent until a recent TDE, it could be a background galaxy with significantly larger associated luminosities and  $M_\bullet$ . Finally, although Maksym et al. (2013) and Donato et al. (2014) both argue that a chance coincidence between the *EUVE* and *Chandra* flare is inherently improbable, the field is sufficiently crowded that within the *EUVE* point spread function (PSF) there exist two brighter, unresolved Sloan Digital Sky Survey (SDSS) objects ( $V < 21$ ; SDSS J134849.21 + 263550.5, SDSS J134850.01+263554.5) of uncertain nature which could conceivably be, e.g. flaring AGNs or Galactic objects.

In order to test these conclusions, we have obtained and analysed deep (8 h) Gemini multislit spectroscopy taken  $\sim 15$  yr after the initial flare, which was necessary in order to confidently determine an absorption-line redshift. In addition, given the excellent quality of the spectrum at redder wavelengths, we model the host galaxy via multicomponent spectral template fitting. We discuss the results of the spectral analysis to determine the classification of the host galaxy of the flare, as well as the two other bright optical objects within the *EUVE* PSF. In addition, we discuss the implications for the flare itself.

Throughout this paper, we adopt concordant cosmological parameters of  $H_0 = 70 \text{ km}^{-1} \text{ s}^{-1} \text{ Mpc}^{-1}$ ,  $\Omega_{\text{m},0} = 0.3$  and  $\Omega_{\lambda,0} = 0.7$ , and calculate distances using Wright (2006). All coordinates are J2000.

## 2 OBSERVATIONS AND DATA

### 2.1 Observations

We observed Abell 1795 from Gemini North with the Gemini Multi-Object Spectrograph (GMOS; Hook et al. 2004) over five nights between 2013 April 8 and 2013 April 19. The queue science observations totalled 8 h and were divided into sixteen 1800 s exposures. We observed at least one Gemini Facility Calibration Unit flat-field per hour of exposure. We used a slit mask with 58 slitlets of varying length, including WINGS J1348, as well as nearby ( $\sim 3.6$  and  $\sim 11.1$  arcsec separation) objects SDSS J134850.01+263554.5 and SDSS J134849.21 + 263550.5 (M13-A1795-S1 and M13-A1795-S2 from Maksym et al. 2013), and four acquisition objects with  $R \sim 17$ –18. The other slitlets covered various other objects which will be discussed elsewhere (Maksym et al., in preparation). The mask was designed from GMOS *r*-band pre-imaging taken in three 3 min exposures on 2013 February 5. Exposures were dithered in position and combined to compensate for chip gaps.

We used the B600 grating, which has an unbinned dispersion of  $0.45 \text{ \AA} \text{ pixel}^{-1}$  and  $2760 \text{ \AA}$  of simultaneous coverage. We took eight exposures per central wavelength at both  $5200$  and  $5250 \text{ \AA}$  in order to compensate for the gaps between the GMOS chips. This configuration afforded wavelength coverage at full exposure between  $3824$  and  $6625 \text{ \AA}$  with rapidly declining sensitivity bluewards of  $4000 \text{ \AA}$ .<sup>1</sup> All slitlets covering science targets had  $0.75$  arcsec widths. GMOS was centred on the coordinates of WINGS J1348, and slitlets were positioned relative to GMOS pre-imaging as determined by *SEXTRACTOR* (Bertin & Arnouts 1996), with WINGS J1348 at  $(\alpha, \delta) = (13^{\text{h}} 48^{\text{m}} 49^{\text{s}}.827, +26^\circ 35' 57''.70)$ , M13-A1795-S1 at  $(\alpha, \delta) = (13^{\text{h}} 48^{\text{m}} 49^{\text{s}}.968, +26^\circ 35' 54''.71)$ , and M13-A1795-S2  $(\alpha, \delta) = (13^{\text{h}} 48^{\text{m}} 49^{\text{s}}.168, +26^\circ 35' 50''.74)$ . Slit lengths and *y*-offsets were  $(7.00; 2.00)$ ,  $(3.70; -1.75)$  and  $(5.00; 0.00)$ , respectively, for optimal spacing. GMOS was oriented  $90$  deg east of north to minimize slit losses at high airmass, since GMOS is not equipped with an atmospheric dispersion corrector. Readout was binned by two pixels in the spatial and spectral directions for an effective pixel size of  $0.145$  arcsec ( $0.91 \text{ \AA}$  per pixel), in order to improve the signal-to-noise ratio (SNR). Conditions were better than 50th percentile for cloud cover and sky brightness, and 70th percentile for seeing (full width at half-maximum, FWHM  $< 0.6$  arcsec). Average airmass was  $\sim 1.1$ , but reached up to  $\sim 1.4$  for science observations.

For absolute photometric calibration, we used recently obtained GMOS spectra of standard star Feige 66 at  $4200$ ,  $5200$  and  $6200 \text{ \AA}$ , and for wavelength calibration we used CuAr arc lamps obtained the same night as the subexposures. From CuAr arc exposures and sky lines of science exposures, we measure  $R \sim 1400$  for this grating configuration.

### 2.2 Data reduction

We processed our data according to the standard GMOS *IRAF* package and tools.<sup>2</sup> We constructed bias files for both science and standard star observations using  $\sim 70$  bias observations obtained by GMOS for each configuration over a  $\sim 2$ -week period encompassing the science observations. We used the flat-field images to automatically edge-detect and cut the spectra. Where the software failed to find subimages (or ‘strips’) associated with individual

<sup>1</sup> [http://www.gemini.edu/sciops/instruments/gmos/gratings/gmos\\_n\\_B600\\_G5303.txt](http://www.gemini.edu/sciops/instruments/gmos/gratings/gmos_n_B600_G5303.txt)

<sup>2</sup> <http://www.gemini.edu/sciops/data-and-results/processing-software>

spectra, we provided manual corrections. We used the `IRAF` tool `GSREDUCE` to remove bias, subtract the overscan values and reject cosmic rays. We calibrated the wavelength scale with `GSWAVELENGTH` and `GSTRANSFORM` by fitting fourth-order Chebyshev polynomials to the CuAr spectra. We subtracted the sky emission by fitting first-order Chebyshev polynomials to the background with `GSSKYSUB`, and we then combined the data into two sets of spectra according to central wavelength by using `GEMCOMBINE`. These spectra were flux-calibrated using `GSCALIBRATE` and sensitivity maps obtained from Feige 66 using `GSSTANDARD`. Given the uncertainties in GMOS absolute photometric calibration for spectra, we folded our spectrum through the SDSS *r*-filter response in order to compare against our pre-imaging and table 3 of Donato et al. (2014). We find our standard star absolute calibration to be  $\sim 0.14$  mag too bright, and we have re-calibrated accordingly.

Using the `PROJECTION` region tool from `DS9`, we extracted spectra from a 0.85 arcsec strip centred on the brightest region. We then combined the exposure-weighted spectra using `IDL`, interpolating the 5200 Å centred spectrum to match the binning of the 5250 Å centred spectrum. Prior to combination, we eliminated instrumental features from chip gaps, as well as poorly subtracted emission from atmospheric [O I] (5578.5 Å, 6301.7 Å) and Na ( $\sim 5890$  Å), substituting the mean continuum values of adjacent areas. We identify one false emission line at  $\sim 4717$  Å which appears to be the result of cosmic rays intersecting over multiple exposures. We measure the SNR as  $F_\lambda/\sigma_{F_\lambda} = (F_{5200} + F_{5250})/\sigma(F_{5200} - F_{5250})$  the standard deviation, where  $F_\lambda$  is the flux for data with grating central wavelength  $\lambda_c$ , and  $\sigma(F_{5200} - F_{5250})$  is the measured standard deviation of the flux difference between data taken with  $\lambda_c = 5200$  and 5250. SNR is  $\sim 2$  near 4000 Å, increases approximately linearly to  $\sim 22$  at 6200 Å and declines again to  $\sim 17$  near 6600 Å. This estimate is comparable to our results from pre-observation simulations using the GMOS-N integration time calculator,<sup>3</sup> which assumed similar conditions.

The final, calibrated spectra, noise curves and residuals are plotted in Figs 1 and 2.

### 3 ANALYSIS

#### 3.1 Pre-imaging

From the guide stars, we use the radial profile fitting function in `ATV` (Barth 2001) to infer a pre-imaging seeing of  $\lesssim 0.67$  arcsec. For WINGS J1348, we measure  $r = 22.16$  (AB), and observed FWHM = 0.78 arcsec. This *r*-value is consistent with values from table 3 of Donato et al. (2014). After correcting for seeing, the FWHM is comparable to the  $\sim 0.3$  arcsec optical FWHM of the dwarf galaxy as determined by Maksym et al. (2013). The observed radial extent is  $\sim 1.0$  arcsec for a circular profile with 25.3 mag arcsec<sup>-2</sup> at  $1\sigma$  above the background. This result is comparable to the result determined by Donato et al. (2014) from the same *Hubble Space Telescope* (*HST*) images analysed by Maksym et al. (2013), and confirms the existence of fainter extended galactic structure. M13-A1795-S1 is unresolved at  $r = 20.52$ , and M13-A1795-S2 is resolved with observed FWHM = 0.84 arcsec and  $r = 20.24$ .

#### 3.2 Absorption-line identification and continuum modelling

We consider the possibility that WINGS J1348 is a background AGN, and initially identify no obvious emission lines in the spectra

by which we may determine a redshift. For a hypothetical background AGN, the  $1\sigma$  limit on an FWHM = 1000 km s<sup>-1</sup> emission line is well described as  $F \lesssim (1.2 \times 10^{-17} \times [\lambda/(4000 \text{ \AA})]^{-7} + 1.6 \times 10^{-18}) \text{ erg cm}^{-2} \text{ s}^{-1}$ . The spectrum appears devoid of strong emission features.

In Maksym et al. (2013), we argued that WINGS J1348 was likely a member of A1795. If so, Ca H+K  $\lambda\lambda 3935, 4214$  Å would be close to 4200 Å, where the SNR remains relatively low. In the case of a dwarf galaxy, the absorption lines might also be intrinsically weak. Having previously used a correlation analysis to investigate whether our 30 min spectrum from the Magellan EchelleTe (Marshall et al. 2008) was consistent with this hypothesis, we do so again with the new, deeper Gemini data. This time, we use `ZFIND` from the `SPEC2D` spectral analysis suite for `IDL`,<sup>4</sup> which cross-correlates the galaxy spectrum with that of a smooth eigenspectrum from the `SPEC2D` template library.

`SPEC2D` finds a family of minima with  $1.09 \lesssim \chi^2/\mu \lesssim 1.12$  at  $0.050 \lesssim z \lesssim 0.067$ , which is comparable to A1795 ( $z = 0.062$ ). It finds a second family of minima with  $1.07 \lesssim \chi^2/\mu \lesssim 1.10$  at  $0.90 \lesssim z \lesssim 0.94$ , and a third family with  $1.08 \lesssim \chi^2/\mu \lesssim 1.13$  at  $1.08 \lesssim z \lesssim 1.11$ .

Since the  $0.050 \lesssim z \lesssim 0.067$  result is consistent with the redshift of A1795, we examine the spectra directly to see if typical small-scale absorption features (not present in the smooth eigenspectrum) correspond to the data near  $z \sim 0.062$ . Ca H+K are difficult to identify unambiguously due to the presence of multiple nearby features of comparable significance. But H $\beta$   $\lambda 4861$  Å and the Mg ib  $\lambda 5173$  Å triplet are prominent with continuum SNR  $\sim 12$ . The Na iD  $\lambda 5893$  Å doublet occurs at continuum SNR  $\sim 22$ . These features require an adjusted  $z = 0.065$ . When fixed at this redshift, other less prominent features become obvious, including Ca H+K, several Fe features, and even the shape of the CN  $\lambda 3883$  Å band (see e.g. Tripicco & Bell 1990). The high quality of the continuum becomes evident when compared to a significantly brighter early-type galaxy (SDSS J134842.56+263700.6,  $M_g = 17.55$ ) from the same observation set, whose continuum closely corresponds to that of WINGS J1348 at scales up to hundreds of angstroms.

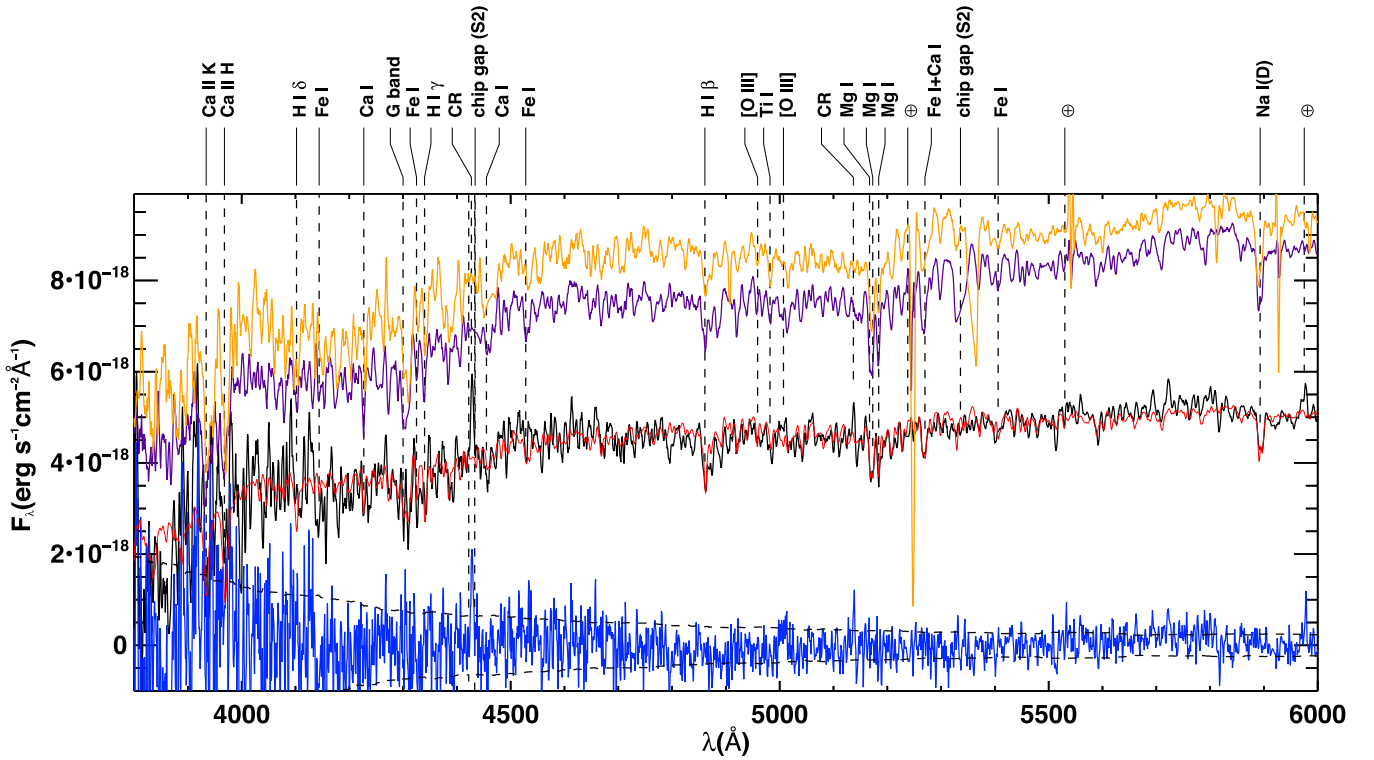
We confirm this analysis via multicomponent fitting of single stellar population models. We have modelled the spectra and measured the absorption-line kinematics using `STARLIGHT` (Cid Fernandes et al. 2005) and Penalized Pixel Fitting (`PPXF`) software (Cappellari & Emsellem 2004). For `STARLIGHT`, we use Bruzual & Charlot (2003) synthetic spectra spanning a grid of 15 ages up to 13 Gyr, and 3 metallicities where  $Z = \{0.004, 0.02, 0.05\}$ , and an expanded grid of 23 ages up to 13 Gyr and  $Z = \{0.0001, 0.0004, 0.004, 0.008, 0.02, 0.05\}$ . Initial `STARLIGHT` models overpredict the continuum by  $\sim 10$  per cent at  $\lambda \lesssim 4600$  Å and  $\lambda \gtrsim 6000$  Å (rest), and require reddening  $A_V \lesssim -0.9$  as per Calzetti et al. (2000). When we apply a polynomial correction to re-calibrate the Gemini data with SDSS J134842.56+263700.6 spectra from SDSS DR10, we find that this issue disappears. We therefore use the re-calibrated data for our modelling.

For `PPXF`, we use Vazdekis et al. (2010) synthetic spectra spanning a grid of 48 ages up to 14 Gyr, with  $[M/H] = \{+0.22, 0.00, -0.40, -0.71, -1.31, -1.71\}$ . These template spectra are convolved with the quadratic difference between the Gemini instrumental resolution and Vazdekis et al. (2010) template resolution. We explore `PPXF` models which include either Calzetti et al. (2000) reddening or which include polynomial corrections to the continuum fit. We find

<sup>3</sup> <http://www.gemini.edu/sciops/instruments/integration-time-calculators/gmosn-itc>

<sup>4</sup> <http://www.astro.princeton.edu/~schlegel/code.html>





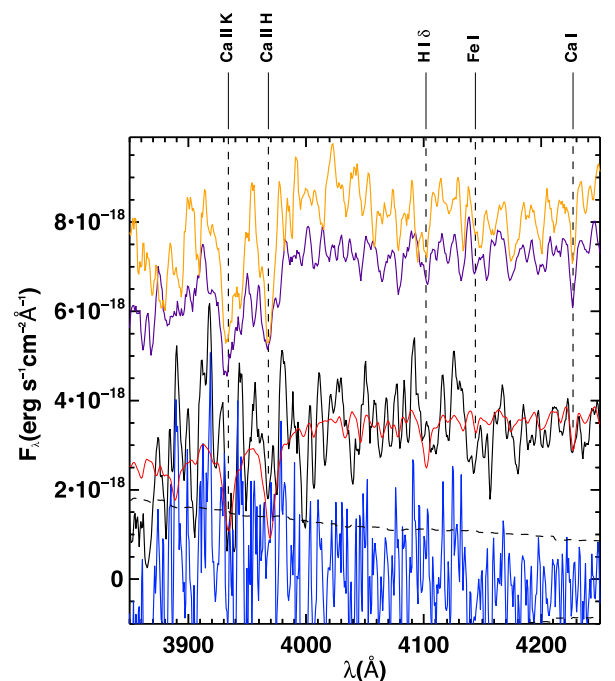
**Figure 1.** Rest-wavelength Gemini spectrum of WINGS J1348, with spectral features of interest indicated. For clarity, the WINGS J1348 data (black, solid) are smoothed with a boxcar average over five  $0.9 \text{ \AA}$  pixels. The best-fitting PPF model (red) which includes Calzetti et al. (2000) internal reddening law is overlaid. Near  $\sim 0 \text{ erg cm}^{-2} \text{ s}^{-1} \text{ \AA}^{-1}$  (bottom: blue, solid) is  $F_\lambda(\text{data}) - F_\lambda(\text{model})$ . Noise (black, horizontal dashed) represents local  $1\sigma$  uncertainty per  $0.9 \text{ \AA}$  pixel, or alternately  $3\sigma$  uncertainty for  $\sim 8 \text{ \AA}$  features. A poorly subtracted cosmic ray (CR) near  $4400 \text{ \AA}$  has been indicated as a noise spike. The atmospheric feature near  $6000 \text{ \AA}$  is [O I]  $\lambda 6363 \text{ \AA}$ . The model is good match for Mg Ib ( $\lambda 5173 \text{ \AA}$ ), Na I D ( $\lambda 5893 \text{ \AA}$ ), the G band ( $\lambda 4300 \text{ \AA}$ ) and numerous stellar absorption lines. Rest-frame spectra from nearby early-type galaxies M13-A1795-S1 (top, orange;  $z = 0.0615$ ) and M13-A1794-S2 (second from top, purple;  $z = 0.0632$ ) have been rescaled to match WINGS J1348, smoothed and overlaid with an offset for comparison. Note that these comparison spectra use  $\lambda_c = 5250 \text{ \AA}$  data only, and thus have not had deep absorption-like features from GMOS chip gaps removed.

that the intrinsic reddening is negligible. The best-fitting PPF model with Calzetti et al. (2000) reddening is overlaid in Figs 1 and 2.

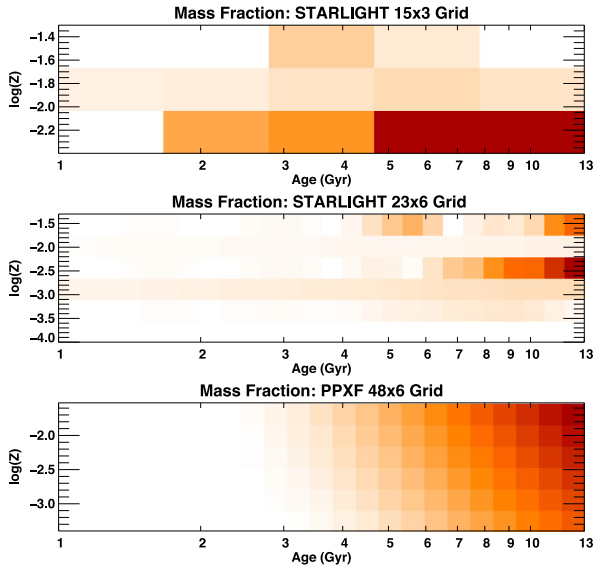
The results of these fits, illustrated in Fig. 3, are consistent with  $\gtrsim 90$  per cent of the stellar mass composed of an old ( $\gtrsim 5 \text{ Gyr}$ ), passive, low-metallicity ( $Z \gtrsim 0.004$ ) population, and possibly also containing a smaller, younger ( $\sim 1 \text{ Gyr}$ ) component with higher metallicity ( $Z \sim 0.05$ ). STARLIGHT also predicts a stellar mass ( $M_*$ ) of only  $\sim (2.3-3.0) \times 10^8 M_\odot$  for the current stellar population. All models find best-fitting radial velocity estimates which are within one spectral dispersion element of  $z = 0.065$ , such that  $z = 0.065 \pm 0.00020$ . Furthermore, all models require velocity dispersions below the instrumental dispersion,  $\sim 50 \text{ km s}^{-1}$  at  $\sim 5500 \text{ \AA}$ . This is consistent with our estimates that the stellar absorption lines are unresolved. Since determinations of velocity dispersions below an instrumental dispersion are unreliable, we take the instrumental dispersion to be an upper limit.

### 3.3 Emission-line measurement

Despite the high quality of the continuum, we do not identify any unambiguous emission features relative to either comparison galaxy spectra in the same field or the PPF best-fitting model. For the purposes of determining the presence of any weak ongoing nuclear or star formation activity, we set upper limits on relevant diagnostic lines covered by our spectrum, [O II]  $\lambda 3727 \text{ \AA}$ , H $\beta$   $\lambda 4861 \text{ \AA}$  and [O III]  $\lambda 5007 \text{ \AA}$ .



**Figure 2.** As Fig. 1, but covering only  $3800-4000 \text{ \AA}$  in order to more clearly display the Ca H+K absorption features. Offsets of M13-A1795-S1 and M13-A1795-S2 have been displaced vertically for clarity.



**Figure 3.** Relative mass fractions of different stellar populations with respect to metallicity and age, as per the best-fitting models from STARLIGHT and PPXF template synthesis codes. Dominant populations are indicated by shading, with darker shading indicating a larger fraction of the total stellar mass in the best-fitting model. Populations with age  $< 1$  Gyr are excluded due to relative lack of significance.

In order to determine the strength of any faint narrow emission lines which are undetected above the continuum, we must assume a line width. The simplest scenario supposes that the velocity of gas in the narrow-line region is comparable to the stellar velocity dispersion ( $\text{FWHM}_{\text{gas}} \sim 0.8 \text{FWHM}_*$ , as per Ho 2009). We therefore assume an upper limit of  $\sigma_{\text{gas}} \lesssim 50 \text{ km s}^{-1}$ , according to the instrumental dispersion and Section 3.2.

We subtract the best-fitting PPXF model with reddening from the measured spectrum and fit the difference to a Gaussian with vertical offset. The inferred  $3\sigma$  upper limits are thus  $F_{[\text{O III}]}$   $< 2.3 \times 10^{-18} \text{ erg cm}^{-2} \text{ s}^{-1}$  or  $L_{[\text{O III}]} < 2.1 \times 10^{37} \text{ erg s}^{-1}$ ,  $F_{\text{H}\beta} < 9.7 \times 10^{-19} \text{ erg cm}^{-2} \text{ s}^{-1}$  or  $L_{\text{H}\beta} < 9.7 \times 10^{36} \text{ erg s}^{-1}$  and  $F_{[\text{O II}]} < 5.4 \times 10^{-18} \text{ erg cm}^{-2} \text{ s}^{-1}$  or  $L_{[\text{O II}]} < 5.0 \times 10^{37} \text{ erg s}^{-1}$ . We expect  $\text{H}\alpha$   $\lambda 6563 \text{ \AA}$  to typically be stronger where Balmer emission is present ( $F(\text{H}\alpha)/F(\text{H}\beta) \sim 3.5$  in AGNs), but  $\lambda 6563 \text{ \AA}$  lies outside our spectral range.

Note that any emission lines from a typical AGN would likely be even narrower (and hence weaker) than this. If, based on the virial theorem, we assume  $M_* \sim 5\sigma_*^2 r_e / G$ , then  $\sigma_* \lesssim 38 \text{ km s}^{-1}$ . Alternately, according to the Faber & Jackson (1976) relation as extended by Jiang, Greene & Ho (2011) to low-mass galaxies, we derive  $\sigma_* \sim 27 \text{ km s}^{-1}$  (assuming  $V - I = 1.34 \text{ mag}$  and  $B - I = 1.8$ , as per Greene, Ho & Barth 2008). There is significant uncertainty to these assumptions, however. In addition to an intrinsic scatter of 0.07 dex to the results of Jiang et al. (2011), their relation is poorly constrained for galaxies so intrinsically faint, with only one galaxy at  $M_I > -16.50$  ( $M_V > -15.16$ ).

These assumptions may be complicated by a possible broad excess observed near  $[\text{O III}] \lambda 5007 \text{ \AA}$ . At only  $\sim 1\sigma$  pixel $^{-1}$ , the SNR of this excess is strongly model dependent ( $\sim 3\sigma - 5\sigma$ ) and from inspection requires  $\text{FWHM} \sim 1000 \text{ km s}^{-1}$  ( $\sim 17 \text{ \AA}$  rest), which is strongly inconsistent with expectations for WINGS J1348 (as per this section). Due to its low significance, we defer further discussion of this putative broad feature to Appendix A. The primary impact of such a broad  $[\text{O III}] \lambda 5007 \text{ \AA}$  feature upon our main

findings would be to suggest a more conservative upper limit for  $[\text{O III}] \lambda 5007 \text{ \AA}$ , such that  $F_{[\text{O III}]} < 1.5 \times 10^{-17} \text{ erg cm}^{-2} \text{ s}^{-1}$  or  $L_{[\text{O III}]} \sim 1.3 \times 10^{38} \text{ erg s}^{-1}$ . But in any event, such emission is unlikely to be caused by illumination of the narrow-line region given the constraints discussed in this section.

### 3.4 Spectroscopy of M13-A1795-S1 and M13-A1795-S2

M13-A1795-S1 and M13-A1795-S objects are  $\sim 1.5 \text{ mag}$  brighter than WINGS J1348 in  $B$ ,  $V$  (Varela et al. 2009) and  $R$  according to the Gemini photometry. As a result, we easily identify their prominent Ca H+K absorption lines, as well as other features typical of early-type galaxies with high confidence. These other objects are passive early-type galaxies common to galaxy clusters, they are members of A1795 ( $z = 0.0615$  and  $0.0632$ , respectively) and they have no evidence of strong line emission typical of AGNs. We have rescaled their rest-frame spectra and plotted them in Figs 1 and 2 for comparison against WINGS J1348.

## 4 DISCUSSION

From these observations, we can immediately infer several important results with respect to Maksym et al. (2013). First and foremost, we unambiguously determine  $z = 0.06514 \pm 0.00020$  (heliocentric) for WINGS J1348. The host galaxy for the *Chandra* 1999–2005 flare described in those papers is therefore clearly a member of A1795 ( $z \sim 0.062$ ). Maksym et al. (2013) had previously supported cluster membership via cross-correlation analysis of a 30 min Magellen Echelle spectrum and association of the galaxy with the cluster’s colour–magnitude ridge line. WINGS J1348 has a moderately high projected relative velocity of  $\sim 900 \text{ km s}^{-1}$ . This is modestly high relative to the cluster dispersion ( $\sim 658 \text{ km s}^{-1}$ ; Cava et al. 2009). Given a projected distance of  $\sim 50 \text{ kpc}$  and a first-order assumption of linear orbits for the cluster galaxies, this relative velocity is thus dynamically consistent with a galaxy near the cluster core. This redshift determination also confirms the distance scales and luminosity distances determined photometrically by Maksym et al. (2013), as well as the inferred properties of the galaxy and associated flare: WINGS J1348 is indeed a passive  $M_V = -14.8$  dwarf galaxy (see also Donato et al. 2014). Via ATV (Barth 2001) profile fitting of the 1999 *HST* F555W observation and the analysis of Donato et al. (2014), we infer a half-light radius of  $r_e \sim 160 \text{ pc}$ .

In addition, the unremarkable early-type spectra of M13-A1795-S1 and M13-A1795-S2 strongly support the hypothesis of Maksym et al. (2013) that the *Chandra* and *EUVE* flare originate from the same event, given that there are no other comparably bright objects in the *EUVE* PSF core (as per Maksym et al. 2013). Neither object is an AGN or Galactic object which could conceivably produce such a bright flare except via another (and comparably rare) X-ray transient.

From both sets of best-fitting STARLIGHT models to WINGS J1348, we infer a stellar mass of  $M_* \sim 2.5 \times 10^8 M_\odot$ , which is somewhat larger than (but comparable to)  $\sim 8.6 \times 10^7 M_\odot$  as derived according to de Jong & Bell (2003) and the  $B - R$  colour (from, e.g. table 3 in Donato et al. 2014). By comparison, Henize 2-10 has  $M_* \sim 3.9 \times 10^9 M_\odot$  (Reines et al. 2011) and the Large Magellanic Cloud has  $M_* \sim 2.7 \times 10^9 M_\odot$  (van der Marel 2006).  $M_*$  for WINGS J1348 is more directly comparable to the least massive galaxies in Reines et al. (2013) or to M60-UCD1 ( $M_* \sim 2 \times 10^8 M_\odot$ ; Strader et al. 2013), though WINGS J1348 is  $> 2$  orders of magnitude less dense. At  $L_V \sim 10^8 L_\odot$  and  $\Sigma \sim 10^{2.7} L_\odot \text{ pc}^{-2}$ , WINGS J1348 appears near the gap between ultra-compact dwarf galaxies and

more luminous dwarf galaxies (fig. 3 in Strader et al. 2013), and given its  $r \gtrsim 600$  pc stellar envelope could represent a transition object, as per Brodie et al. (2011). At a basic level, template fitting reveals a very old stellar population of uncertain or mixed metallicity, as suggested by Maksym et al. (2013). The star formation rate inferred from our limit to  $H\beta$  emission as per Moustakas, Kennicutt & Tremonti (2006) is negligible at  $\lesssim 4.4 \times 10^{-4} M_{\odot} \text{ yr}^{-1}$  and strongly reduces the likelihood of a core-collapse supernova explanation for the flare. With so little ongoing star formation, WINGS J1348 is unlikely to be a dwarf irregular and therefore may be a dwarf spheroidal. But without better structural and kinematical information, any more precise label than ‘dwarf galaxy’ cannot be applied with certainty.

An assumed range of  $10^8 \lesssim M_{\bullet} \lesssim 3 \times 10^8 M_{\odot}$  allows us to derive an expected  $M_{\bullet}$ , assuming that the stellar mass resides entirely in a spheroid. In this case, the  $M_{\text{bulge}} - M_{\bullet}$  relationship of Kormendy & Ho (2013) implies  $\log(M_{\bullet}/M_{\odot}) \sim 5.1 - 5.7$ , which is consistent with Maksym et al. (2013) and Donato et al. (2014). This range is also consistent with an upper limit on  $M_{\bullet}$  based on the best-fitting kinematics of the absorption-line spectrum, which imply an unresolved, lower mass black hole. Given  $\sigma_{\star} \lesssim 50 \text{ km s}^{-1}$ , we expect  $M_{\bullet} \lesssim 7 \times 10^5 M_{\odot}$  (Kormendy & Ho 2013). With greater confidence in the *EUVE* flare’s TDE origin, we also infer  $M_{\bullet} \gtrsim 2 \times 10^5 M_{\odot}$ , assuming that the TDE is Eddington-limited.

We strongly constrain any persistent nuclear activity from WINGS J1348, and improve upon previous limits from Maksym et al. (2013) and Donato et al. (2014). In particular, we set a limit on narrow  $H\beta$  a factor of  $\sim 180$  lower than Maksym et al. (2013). Even with our most conservative upper limits on weak [O III] emission, we also infer that any nuclear activity must be temporary or very weak. By comparison, Panessa et al. (2006) and Ho, Kim & Terashima (2012, fig. 3) imply  $L_{\text{X}}(0.5 - 2.0 \text{ keV}) \sim 1.5 \times 10^{39} \text{ erg s}^{-1}$ , or a factor of  $\sim 2 \times 10^4$  below the 1998 *EUVE* flux, and  $\sim 20$  below the 2010 *Chandra* upper limit in Maksym et al. (2013). We derive comparable results from Stern & Laor (2012),  $L_{\text{X}} \sim 2.6 \times 10^{39} \text{ erg s}^{-1}$  and  $L_{\text{bol}} \sim 6.9 \times 10^{39} \text{ erg s}^{-1}$ . The inferred low-level AGN emission is reduced even further if we assume narrower [O III] widths as per Ho (2009) and Section 3.3, such that  $L_{\text{X}}$  and  $L_{\text{bol}}$  are reduced by factors of  $\sim 6$  and 13, respectively. Our spectrum does not cover the useful BPT (Baldwin, Phillips & Terlevich 1981) lines redwards of 6300 Å, but we infer that the  $H\beta$  is consistent with the Maksym et al. (2013) non-detection of  $H\alpha$  assuming  $F(H\alpha)/F(H\beta) \sim 3.5$ .

## 5 CONCLUSIONS

Via new deep Gemini spectroscopy we have obtained a redshift for the faint ( $V \sim 22.5$ ) host of the X-ray flare identified in archival *Chandra* data by Maksym et al. (2013) and later by Donato et al. (2014), and have conducted a detailed analysis of its spectral properties. We strongly confirm the major assertions of Maksym et al. (2013) and Donato et al. (2014), namely the following.

- (i) Our derived redshift agrees with the initial redshift estimate from Maksym et al. (2013). The flare therefore occurred in a dwarf galaxy in Abell 1795, with only  $M_V \sim -14.8$ . The host galaxy, WINGS J1348, has a half-light radius of  $r_e \sim 160$  pc and an extended stellar envelope of  $\gtrsim 600$  pc.
- (ii) All luminosities from Maksym et al. (2013) and Donato et al. (2014) based upon an assumed redshift of  $z \sim 0.062$  are also supported.
- (iii) The luminous ( $\gtrsim 3 \times 10^{43} \text{ erg s}^{-1}$ ) *EUVE* flare of 1998 March is likely to have the same origin as the *Chandra* flare, given that the

two brighter objects within  $\sim 15$  arcsec of WINGS J1348 are both non-active early-type members of A1795.

- (iv) A tidal flare explanation is strongly favoured, given limits on AGN emission which suggest no more than weak ( $\lesssim 10^{39} \text{ erg s}^{-1}$ ) ongoing accretion, as well as a passive early-type galaxy spectrum with marginal star formation and an old ( $\gtrsim 5$  Gyr) stellar population.
- (v) The host galaxy is therefore one of the smallest galaxies confirmed to host an MBH, and is likely to host an IMBH. Limits from  $M_{\star}$  ( $\log[M_{\star}/M_{\odot}] \sim 5.1 - 5.7$ ) are consistent with those from Eddington and stellar dispersion limits ( $\log[M_{\star}/M_{\odot}] \sim 5.3 - 5.8$ ).
- (vi) The TDE rate therefore may be favourably affected by the presence of IMBHs in cluster dwarf galaxies, as per Wang & Merritt (2004). A larger data set, however, is necessary to strongly constrain the rate from such objects.

In addition, we find the following.

- (i) The implied total decline from the peak X-ray flux is a factor of  $\gtrsim 10^4$  over a period of  $\sim 15$  yr.
- (ii) WINGS J1348 is moving at  $\sim 900 \text{ km s}^{-1}$  relative to A1795, dynamically consistent with a linear orbit in the cluster core at a projected distance of  $\sim 50$  kpc.
- (iii) The MBH is unlikely to have been ejected from an MBH–MBH merger given its low stellar dispersion ( $\lesssim 50 \text{ km s}^{-1}$ ) and (as per Maksym et al. 2013) large stellar mass indicative of a dwarf galaxy.

These results support the presence of IMBHs in normally inactive dwarf galaxies, and complement efforts to understand the MBH population of low-mass galaxies such as via observations of AGNs (Reines et al. 2013) and weakly accreting black holes (Miller et al. 2012).

Potential opportunities for future investigation include deep spectroscopy of WINGS J1348 at  $\gtrsim 6300$  Å to better constrain emission from [O I],  $H\alpha$ , [N II] and [S II] diagnostic lines. Also, deep observations at higher spectral resolution could better determine  $\sigma_{\star}$ , and hence directly compare the IMBH of this extremely low mass galaxy with the  $M_{\bullet} - \sigma$  relation. Such measurements would be of interest for comparison against measured environmental effects in the  $M_{\bullet}$  relations, such as by McGee (2013).

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## APPENDIX A: A POSSIBLE BROAD EXCESS NEAR O III $\lambda 5007 \text{ \AA}$

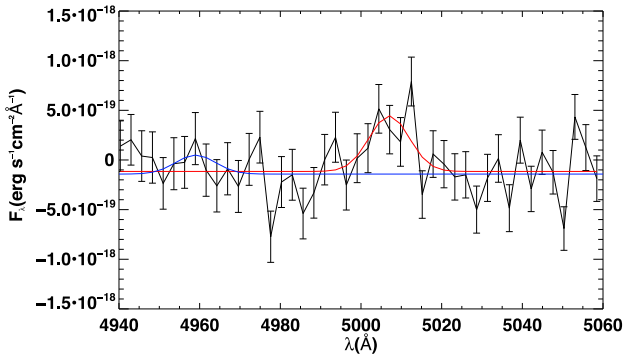
In Section 3.3, we mentioned a possible broad excess within  $\sim 9 \text{ \AA}$  of O III  $\lambda 5007 \text{ \AA}$ . We have noted that at only  $\sim 1\sigma$  pixel $^{-1}$  such a broad flux excess spans FWHM  $\sim 1000 \text{ km s}^{-1}$  if it is taken to be statistically significant, which is much greater than measured  $\sigma_* \lesssim 50 \text{ km s}^{-1}$  and therefore incompatible with standard assumptions regarding narrow-line emission from AGNs (see e.g. Ho 2009).

Emission from [O III] is, however, predicted as long-term post-disruption feature by various TDE models (e.g. Eracleous, Livio & Binette 1995; Clausen & Eracleous 2011; Clausen et al. 2012) and has been associated with TDE candidates (Irwin et al. 2010; Cenko 2011). Such emission may be asymmetrical and comparably broad, owing to illumination of the tidal debris stream and ejecta, which continue to evolve at over long time-scales and occupy spatial scales much smaller than the narrow-line region of AGNs. The detection of such features would therefore be of interest in comparison with such models. We therefore describe it here for the sake of completeness.

In Fig. A1, we display the region within  $50 \text{ \AA}$  of  $\lambda 5007 \text{ \AA}$  (rest-frame), re-binned to  $2.7 \text{ \AA}$  width for clarity. The statistical significance is insufficient to determine a width via Gaussian fitting, so we estimate the FWHM from inspection of the unbinned data, with an uncertainty no less than  $\sim 200 \text{ km s}^{-1}$ . Without accounting for systematic uncertainties in the continuum model, which are beyond the scope of this paper and may be significant, we find the putative detection to be  $> 3\sigma$  when fitting a simple Gaussian to the unbinned data within  $100 \text{ \AA}$  of  $\lambda 5007 \text{ \AA}$ , whether using the best-fitting PXP model or by rescaling the early-type galaxy SDSS J134842.56 + 263700.6 from the same data set to match the continuum. When the Gaussian model includes a vertical offset to match the PXP continuum,  $\sigma = 4.8$ . We obtain comparable results when directly summing the flux excess.

Given possible systematic errors, the putative detection is therefore marginal and uncertain. Also, any [O III]  $\lambda 5007 \text{ \AA}$  should be accompanied by [O III]  $\lambda 4959 \text{ \AA}$  with  $\sim 1/3$  of the  $\lambda 5007 \text{ \AA}$  flux. Such confirmation is not possible given the low SNR of [O III]





**Figure A1.**  $\lambda 5007 \text{ \AA}$  region (rest-frame), re-binned to  $2.7 \text{ \AA}$  width for clarity. Red: best-fitting Gaussian with vertical offset for a putative [O III] line centred at  $\lambda 5007 \text{ \AA}$ . Blue: the same Gaussian, centred at  $\lambda 4959 \text{ \AA}$  and rescaled by  $1/3$ .

$\lambda 5007 \text{ \AA}$  (see Fig. A1); however, any such  $\lambda 4959 \text{ \AA}$  flux could be hidden by a  $\gtrsim 1\sigma$  fluctuation. There is no significant broad Balmer emission with the same FWHM, most notably for  $H\beta$   $\lambda 4861 \text{ \AA}$  ( $F_{H\beta} < 5.5 \times 10^{-18} \text{ erg cm}^{-2} \text{ s}^{-1}$ ). Since common emission lines are masked in the PPF fit, such  $H\beta$  is unlikely to be hidden by filling in the absorption line.

Considering the possibility that the excess could be spurious, we examined the 16 subexposures individually and found no evidence for cosmic rays or instrumental artefacts. We also compare against the co-added spectra from both  $\lambda 5200$  and  $\lambda 5200 \text{ \AA}$  central wavelengths. Each subset of eight exposures has an excess of  $\sim 1.3 \times 10^{-17} \text{ erg cm}^{-2} \text{ s}^{-1}$  within  $12 \text{ \AA}$  of  $5007 \text{ \AA}$ , comparable to the best-fitting value for the full data set.

A detailed treatment of the emission mechanics of the disruption is beyond the scope of this paper, but we note some basic constraints. Assuming a cloud-and-stream model simplified from

Guillochon, Manukian & Ramirez-Ruiz (2014) to be stellar debris and a solar-type star disrupted at the tidal radius by an IMBH with  $M_{\bullet} \sim 10^{5.4} M_{\odot}$ , after  $\sim 15 \text{ yr}$  the marginally bound material will have expanded to  $\sim 2 \times 10^{14} \text{ cm}$ , less than  $r \sim 3.4 \times 10^{15} \text{ cm}$  implied for Keplerian motion. The area of the bound stream may be significant ( $\sim 2 \times 10^3 \text{ au}^2$ ), with a range of velocities  $\lesssim 9 \times 10^3 \text{ km s}^{-1}$ , and may include an outer sheath reaching the low critical density of [O III]  $n \sim 6 \times 10^5 \text{ cm}^{-3}$ . The data are insufficient, however, to demonstrate significant line-of-sight velocity which is likely from such a stream (Strubbe & Quataert 2009).

Several interpretations of the putative [O III] line are possible which do not necessitate gas dominated by the IMBH gravity or the dynamics of the debris cloud and ejecta stream. For example, a radiatively driven wind could occur at large ( $\lesssim 4.6 \text{ pc}$ ) radii given an initial impulse from the early  $\gtrsim 10^{43} \text{ erg s}^{-1}$  flare. Such winds have been attributed to [O III] emission from the Virgo elliptical galaxy NGC 4472 globular cluster RZ 2109, possibly in conjunction with a stellar-mass black hole, and require relatively little gas mass ( $\lesssim 10^{-3} M_{\odot}$  for  $L_{\lambda 5007} \sim 10^{37} \text{ erg s}^{-1}$ ; Steele et al. 2011). Such gas likely exceeds the escape velocity of the galaxy. Note also that if the galaxy is moving through the cluster core, it is likely to be stripped of most of its gas, and what gas remains may be strongly influenced by the intracluster medium (ICM). For comparison,  $1000 \text{ km s}^{-1}$  is comparable to the line-of-sight motion of WINGS J1348 relative to A1795. Also, the temperature-derived ICM sound speed is  $\sim (5kT/3\mu m_p)^{1/2} \sim 700\text{--}1300 \text{ km s}^{-1}$  for  $kT \sim 2\text{--}6 \text{ keV}$  plasma (as measured by Gu et al. 2012), implying that the motion of WINGS J1348 could be supersonic and might thus shock any gas still retained by the galaxy.

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